

# Extraordinary transmission, symmetry and the Blaschke Product

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## Extraordinary optical transmission

Extraordinary optical transmission (EOT) is the phenomenon of greatly enhanced transmission of light. It goes back to Ebbesen, et. al. "Extraordinary optical transmission through sub-wavelength hole arrays," *Nature* 391 (1998), 667. It is thought to be caused by [surface plasmons](#).

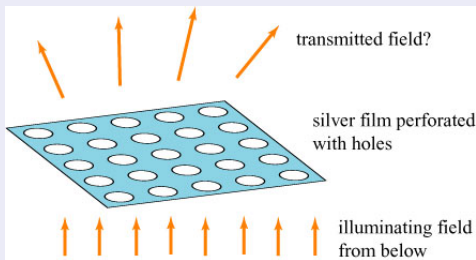


Figure: Schematic of Extraordinary optical transmission

# Extraordinary acoustic transmission

Extraordinary acoustic transmission is the acoustic analogy of extraordinary optical transmission.

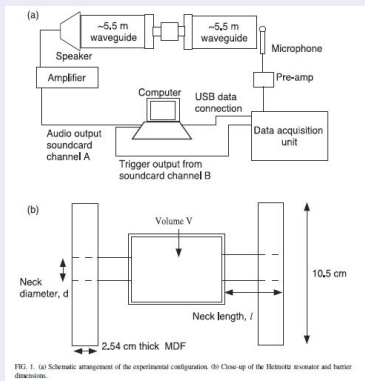


Figure: Schematic of [Extraordinary acoustic transmission](#) from Crow et. al. *AIP Advances* 5, 027114 (2015)

This figure from Crow et. al. shows the **key results** of the experiment

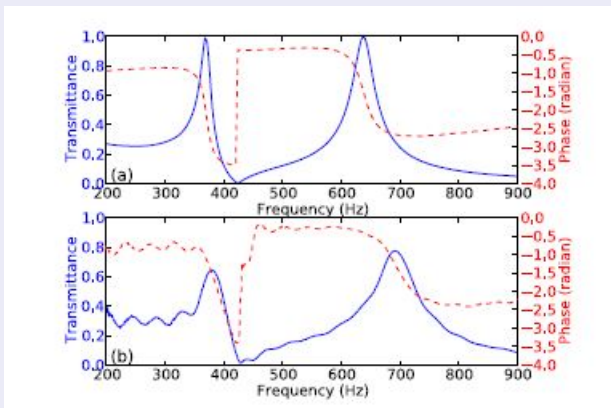


Figure: Figure 6 from Crow et. al.

## Symmetry

We want to explain this extraordinary transmission and to give conditions for it to exist. **Symmetry** plays a key role. We assume we have **one mode** of propagation and that the problem is symmetric.

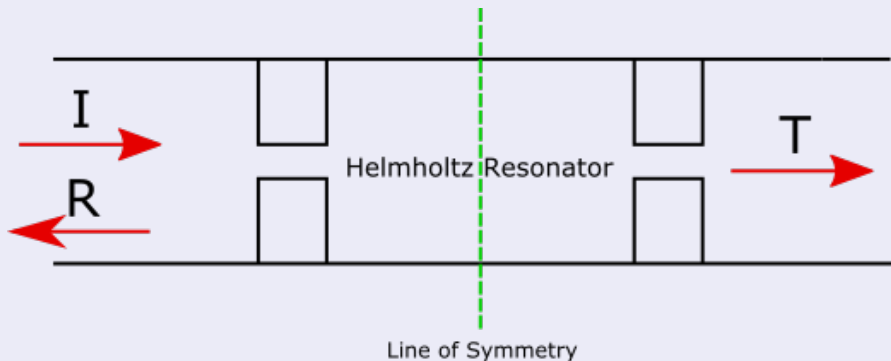


Figure: Schematic diagram of our problem

We can decompose the problem into the following two subproblems.

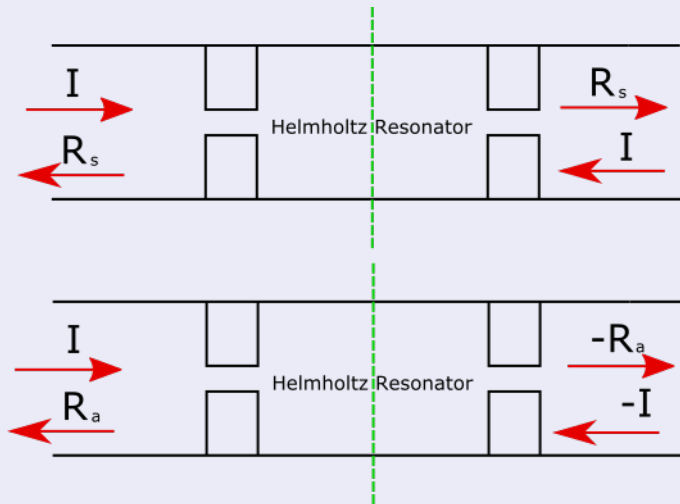


Figure: Symmetric and antisymmetric problem

If we compare the three problems

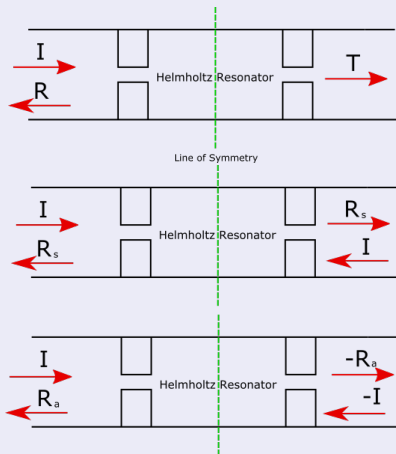


Figure: It follows that  $R = \frac{1}{2} (R_s + R_a)$  and  $T = \frac{1}{2} (R_s - R_a)$

## Decomposition of the Reflection

The reflection coefficients  $R_s$  and  $R_a$  are **analytic functions** of the frequency  $k$ . For a wave guide they are **meromorphic** with poles in the lower half plane.

We know that, from conservation of energy

$$|R_s(k)| = |R_a(k)| = 1, k \in \mathbb{R}$$

We can then conclude that  $R_s$  is basically a **Blaschke Product**

$$R_s(k) = e^{if(k)} \prod_i \left( \frac{k - \bar{a}_m}{k - a_m} \right)$$

where  $f(k)$  takes real values for real  $k$ .



# Blaschke Product

$$R_s(k) = e^{if(k)} \prod_i \left( \frac{k - \bar{a}_m}{k - a_m} \right)$$

- The values  $a_m$  are the points where the analytic extension of the scattering matrix is singular .
- The  $a_m$  are known as **resonances**.
- They are the points where the resolvent is singular and they are **almost** eigenvalues.
- Without symmetry the scattering matrix has dimension two and we cannot apply this decomposition.

# Resonances and Helmholtz Resonators

- A Helmholtz resonator is a **physical object** - a bottle which you blow over for example.
- Mathematically we have **perturbed** the eigenvalues which have become resonances.
- Resonances and eigenvalues are **not** the same thing. For example the mode associated with a resonances **grows** with distance away.

# From the Blaschke Product to Extraordinary Transmission

The form of the transmission is

$$T = \frac{1}{2} (R_s - R_a)$$

and we know that

$$|R_s(k)| = |R_a(k)| = 1, k \in \mathbb{R}$$

and

$$R_s(k) = e^{if(k)} \prod_i \left( \frac{k - \bar{a}_m}{k - a_m} \right)$$

Therefore near  $a_m$  there is a **change of phase of  $2\pi$** . This in turn implies somewhere near  $a_m$ ,  $|T| = 1$ . We are assuming that  $f(k)$  is slowly varying.

## Example Calculations

The governing equation is

$$\nabla^2 \phi(x, y) + k^2 \phi(x, y) = 0$$

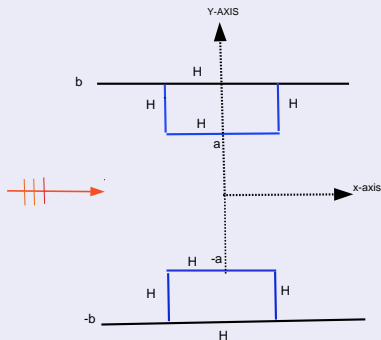


Figure: Our geometry

## Example Calculations

The solutions is found by **mode matching**

$$\phi(x, y) = \sum_{n=1}^{\infty} A_n^{(s)} e^{-i\bar{\alpha}_n(x+l)} \psi_n(y) + e^{i\bar{\alpha}_1(x+l)} \psi_1(y)$$

where

$$\alpha_n = \frac{(n-1)\pi}{2b}, \quad \bar{\alpha}_n = \sqrt{(k^2 - \alpha_n^2)},$$

and

$$\psi_n(y) = \begin{cases} \sqrt{\frac{1}{b}} \cos \alpha_n(y-b), & n \neq 1, \\ \sqrt{\frac{1}{2b}}, & n = 1, \end{cases}$$

# Analytic Extension

There are two challenges to this solution method.

- 1 How to find the **analytic extension** of the solution for complex  $k$ ?
- 2 How to find the **singularities**  $a_m$ ?

These are overcome by

- 1 A **homotopy method**. The roots should not **jump across branch cuts**.
- 2 **Newton's method** combined with a complex variable **bisection method**.

# Results

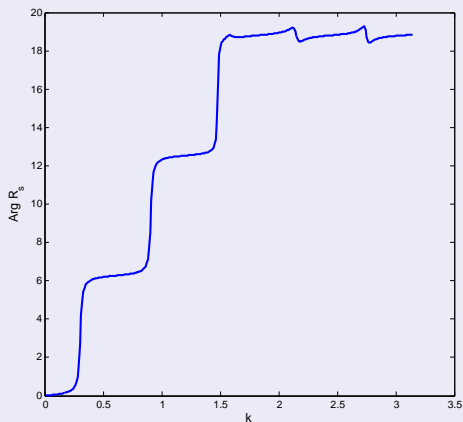


Figure: Reflection coefficient against  $\text{Arg}(R_s)$

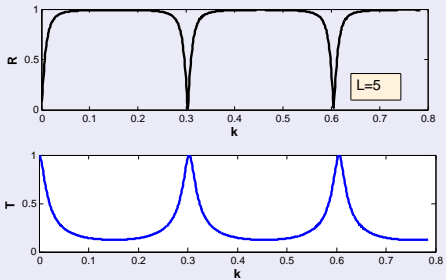


Figure: Reflection and Transmission against wave number for  $L = 5$



# Conclusions

- If there is a **symmetry** and a **single mode of propagation** we will get **extraordinary transmission**, for each resonance.
- We have tested this theory for a simple case in acoustic scattering similar to a recent experimental verification.
- There are many interesting questions raised by this research, ranging from **proof of analyticity** to possible applications in wave splitting or even an **acoustic prism**.